#### **III BRUNEL Workshop on Random Matrix Theory**

# Matrices, Characters, Quantum (Super)Spin Chains

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Brunel University, Uxbridge, 17/12/07

with P.Vieira, arXiv:0711.2470 with A.Sorin and A.Zabrodin, hep-th/0703147.

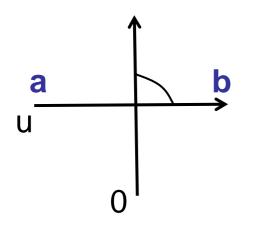
### **Motivation and Plan**

- Classical and quantum integrability are intimately related (not only through classical limit!). Quantization = discretization.
- Quantum spin chain Discrete classical Hirota dynamics for fusion of quantum states (according to representation theory)
   Klumper, Pearce 92', Kuniba, Nakanishi, '92, Krichever, Lupan, Wiegmann, Zabrodin'97
- Based on Bazhanov-Reshetikhin (BR) formula for fusion. Bazhanov, Reshetikhin'90
   Direct proof of BR formula seems to be absent (but see Cherednik'88)
   We fill this gap using the gl(K|M) character technique. v.k., Vieira'07
- Solution of Hirota eq. for (super)spin chain in terms of Baxter TQ-relation.

Tsuboi'98 V.K., Sorin,Zabrodin'07

- More general and more transparent with SUSY: new QQ relations.
- Alternative to algebraic Bethe ansatz:
   all the way from R-matrix to nested Bethe Ansatz Equations
- Possible random matrix applications: from characters to matrix integrals.

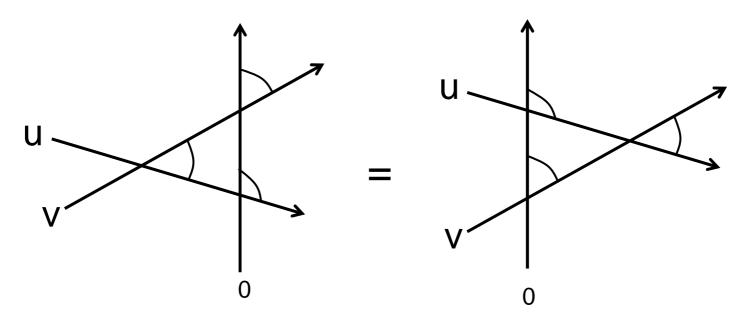
# gl(K|M) super R-matrix and Yang-Baxter



$$R(u) = u + 2\sum_{\alpha\beta} (-)^{p_{\alpha}} e_{\beta\alpha} \otimes e_{\alpha\beta}$$
  
=  $u + 2\mathcal{P}$  (super)permutation

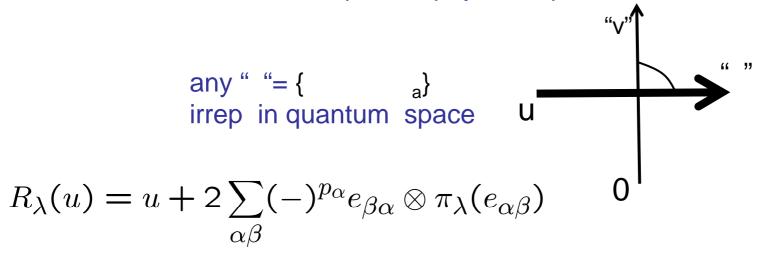
$$p_{\alpha} = 0(1);$$
 for Desen(Fermion) indexes

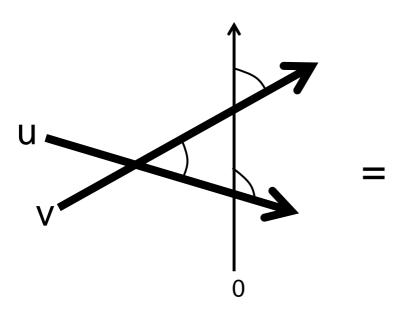
GI(K|M) generator: 
$$\left(e_{\alpha\beta}\right)^{ab}=\delta_{a\alpha}\delta_{b\beta}$$



# Fused R-matrix in any irrep of gl(K|M)

vector irrep "v" in physical space





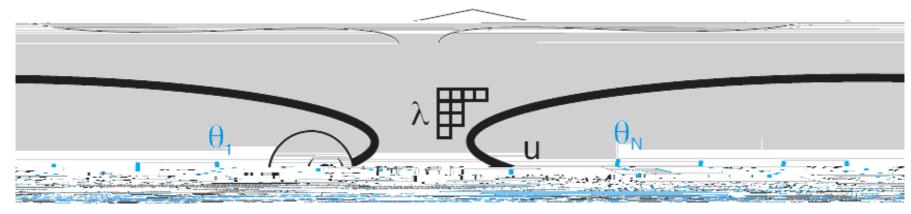
### **Twisted Transfer Matrix**

$$T^{\{\lambda\}}(u) = \operatorname{str}\left(\mathcal{L}^{\{\lambda\}}(u)\pi_{\lambda}(g)\right) \qquad \begin{array}{c} \operatorname{polynomial} \\ \operatorname{of\ degree\ N} \end{array}$$

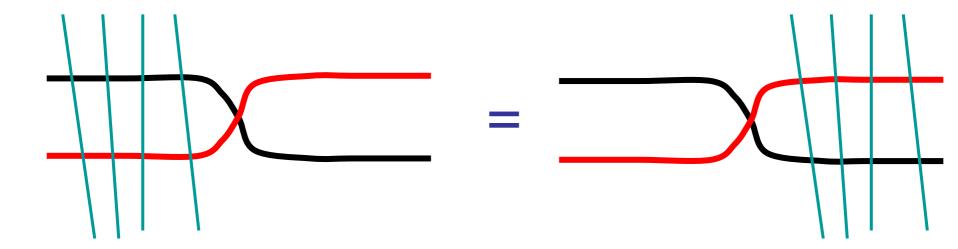
• Defines all conserved charges of (inhomogeneous) super spin chain:

$$\left[T^{\{\lambda\}}(u),T^{\{\lambda'\}}(u')\right]=0$$

$$T_{\lambda}(u)$$



### Commutativity of T-matrices using Yang-Baxter



$$L(u) L(v) R(u-v) = R(v-u) L(v) L(u)$$

#### Bazhanov-Reshetikhin fusion formula

$$T^{\{\lambda\}}(u) = \frac{1}{S_N(u)} \det_{1 \le i,j \le a} T_{\lambda_i + i - j}(u + 2 - 2i)$$

Bazhanov, Reshetikhin'90 Cherednik'87

$$S_N(u) = \prod_{n=1}^N \prod_{k=1}^{a-1} (u - \theta_a - 2k)$$

• Expresses  $T^{\{\lambda\}}(u)$  for general irrep ={ 1, 2,..., a} through  $T_s(u)$  in symmetric irreps

Compare to Jacobi-Trudi formula for GL(K|M) characters

$$\chi_{\{\lambda\}}(g) = \det_{1 \le i,j \le a} \chi_{\lambda_j + i - j}(g).$$

 $\chi_s(g)$  - symmetric (super)Schur polynomials with generating function

$$\vec{w}(z) \equiv \operatorname{sdet}(1-zg) \stackrel{1}{=} \sum_{s=1}^{\infty} \chi_s(\vec{g}) z^s, \qquad g \in gl(K|M)$$

#### T-matrix and BR formula in terms of left co-derivative

Monodromy matrix:

$$\hat{L}^{\{\lambda\}}(\underline{u}) = (u_1 + 2\hat{D}) \otimes (u_2 + 2\hat{D}) \otimes \ldots \otimes (u_N + 2\hat{D}) \pi_{\lambda}(g), \qquad u_n \equiv u - \theta_n$$

- Transfer-matrix of chain without spins:  $T^{\{\lambda\}} = \text{str } \pi_{\lambda}(g) = \chi_{\{\lambda\}}(g)$
- Transfer-matrix of one spin:  $T^{\{\lambda\}}(\underline{u}) = (\underline{u} + 2\widehat{D}) \chi_{f\lambda\lambda}(\underline{q})$

Transfer-matrix of N spins

$$T^{\{\lambda\}}(u) = (u_1 + 2\widehat{D}) \otimes (u_2 + 2\widehat{D}) \otimes \ldots \otimes (u_N + 2\widehat{D}) \chi_{\{\lambda\}}(g)$$

#### Proof of BR formula for one spin

Jacobi-Trudi formula for character

#### should be equal to

$$T_{BR}^{\{\lambda\}}(u) \equiv \det_{1 \le i,j \le a} \left[ (u+2-2i+2\widehat{D})\chi_{\lambda_i+i-j}(g) \right]$$

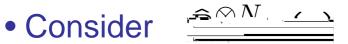
• First, check for trivial zeroes: every 2x2 minor of two rows

$$T_{BR}^{\{\lambda\}}(2k) = \begin{pmatrix} \dots & T_{\lambda_k+k-1}(2) & T_{\lambda_{k+1}+k}(0) & \dots \\ \dots & T_{\lambda_k+k-2}(2) & T_{\lambda_{k+1}+k-1}(0) & \dots \\ \dots & \dots & \dots & \dots \\ \dots & T_{\lambda_k+k-a}(2) & T_{\lambda_{k+1}+k-a+1}(0) & \dots \end{pmatrix}$$

is zero due to curious identity for symmetric characters

$$(1 + \hat{D})_{x_{s_1}} = \hat{D}_{x_{s_2}} = (1 + \hat{D})_{x_{s_1} + 1} = \hat{D}_{x_{s_2} + 1} = 0$$

# Proof of the main identity



One derivative

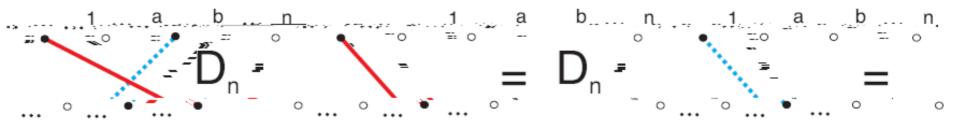
$$\widehat{D}w(z) = \frac{gz}{1 - gz}w(z) \Leftrightarrow \left[\widehat{D}w(z)\right]_{j_1}^{i_1} = \left[\frac{gz}{1 - gz}\right]_{j_1}^{i_1}w(z)$$

$$D w(z) = w(z)$$

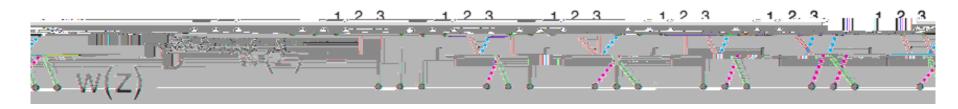
### Graphical representation

$$\widehat{\mathbb{D}} \otimes \frac{gz}{1-g\overline{z}} = \mathcal{P}\left(rac{1}{1-gz} \otimes rac{gz}{1-gz}
ight)$$

One derivative



Three derivatives



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### Proof of the main identity (continuation)

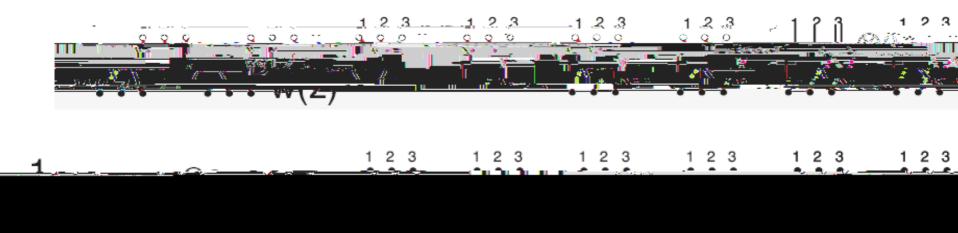
Consider

$$(1+\hat{D})^{\otimes N}w(z)$$

$$(1+\hat{D})w(z) = \frac{1}{1-gz}w(z) \Leftrightarrow \left[ (1+\hat{D})w(z) \right]_{j_1}^{i_1} = \left( \frac{1}{1-gz} \right)_{j_1}^{i_1} w(z)$$

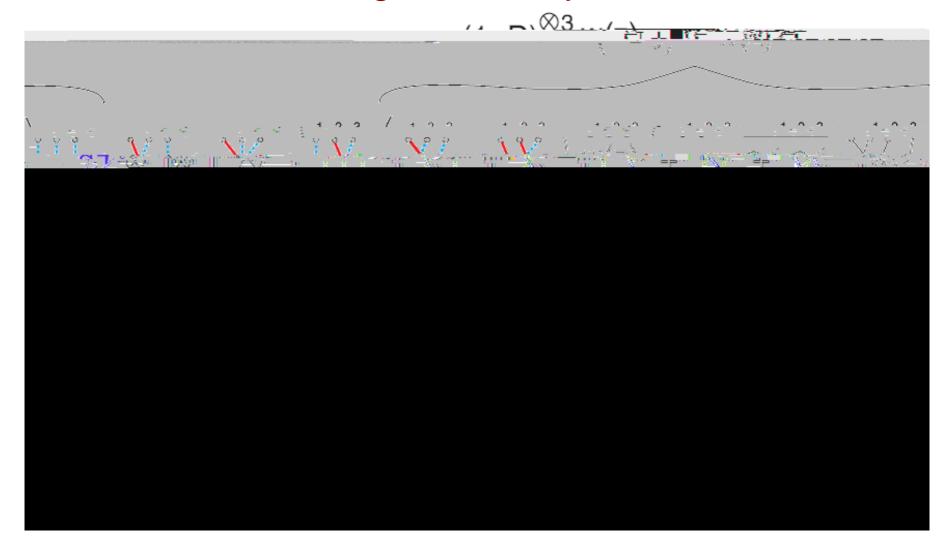
• In pictures:

## Comparison



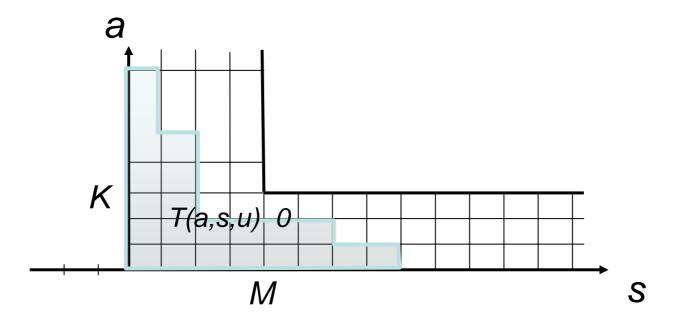
- Notice that the difference is only in color of vertical lines.
- Identical after cyclical shift of upper indices to the right in 2-nd line (up to one line where red should be changed to dotted one)

### Proving the identity ....



• This completes our proof of Bazhanov-Reshetikhin formula

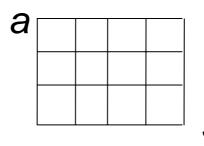
### SUSY Boundary Conditions: Fat Hook



All super Young tableaux of gl(K|M) live within this fat hook

### Hirota eq. from Jacobi relation for rectangular tableaux

T(a,s,u):



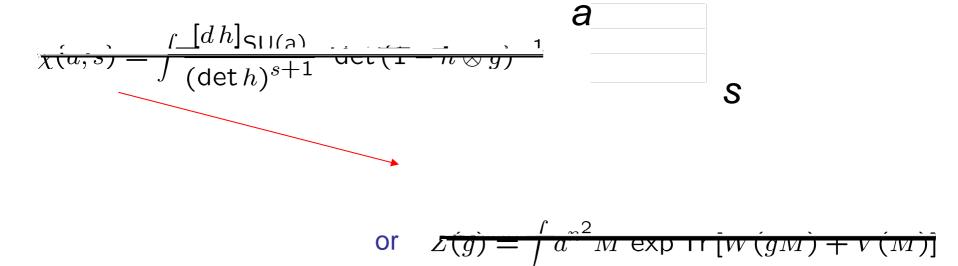
From BR formula, by Jacobi relation for det:

we get Hirota eq.

$$T(a, s, u+1)T(a, s, u-1)-T(a, s+1, u)T(a, s-1, u) = T(a+1, s, u)T(a-1, s, u)$$

We can solve it by Bäcklund trick, find

# From characters to matrix integrals



or any invariant function of 
$$g \in GL(n)$$

• Example: The matrices acting in the space  $(\mathbf{V}_n)^{\otimes N}$ 

$$T^{\{\lambda\}}(u) = (u + u_1 + 2\hat{D}_g) \otimes (u + u_2 + 2\hat{D}_g) \otimes \ldots \otimes (u + u_N + 2\hat{D}_g) \langle \operatorname{Tr} \frac{1}{1 - gM} \rangle$$

Commute for any 
$$u$$
:  $\left[T(u), T(u')\right] = \Im$ 

Interesting relations for random matrix correlators!

#### Solution: Generalized Baxter's T-Q Relations

[V.K.,Sorin,Zabrodin'07]

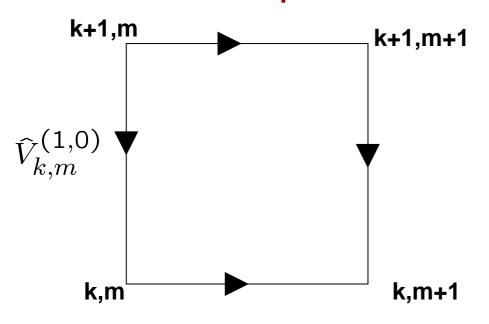
• Diff. operator generating all T's for symmetric irreps:

Introduce shift operators:

$$\frac{\widehat{V}(1,0)_{\ell_{a},N} - Q_{k+1,m}(u) Q_{k,m}(u+2)}{Q_{k+1,m}(u+2) Q_{k,m}(u)} \qquad \qquad \downarrow \\
\widehat{V}_{k,m}^{(0,-1)}(u) = y_{m} \frac{Q_{k,m}(u) Q_{k,m+1}(u+2)}{Q_{k,m}(u) Q_{k,m+1}(u+2)} - e^{2\partial_{u}} \qquad \qquad \downarrow$$

Baxter's Q-functions:

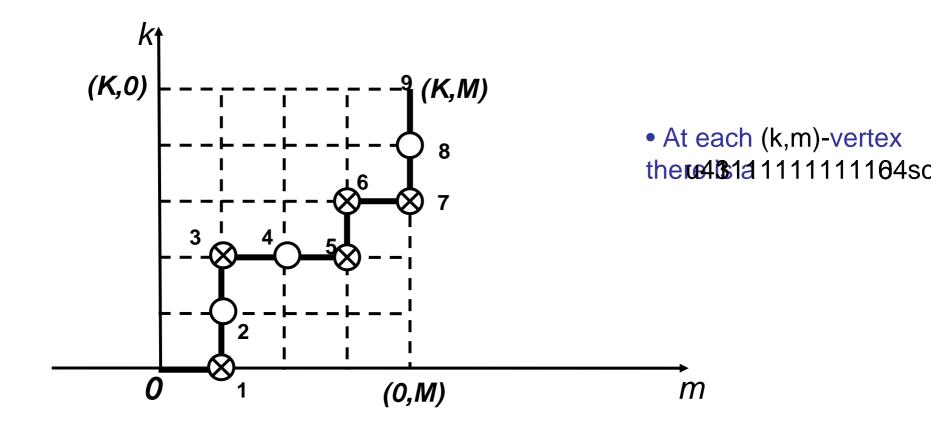
### Hirota eq. for Baxter's Q-functions



Zero curvature cond. for shift operators

$$Q_{k,m}(u)Q_{k+1,m+1}(u+2) - Q_{k+1,m+1}(u)Q_{k,m}(u+2) = Q_{k,m+1}(u)Q_{k+1,m}(u+2)$$

### Undressing along a zigzag path (Kac-Dynkin diagram)



### Bethe Ansatz Equations along a zigzag path

BAE's follow from zeroes of various terms in Hirota QQ relation

$$\prod_{b=1}^{K+M} \frac{\breve{Q}_b \left( \breve{u}_j^{(a)} - K_{ab} \right)}{\breve{Q}_b \left( \breve{u}_j^{(a)} + K_{ab} \right)} = (-1)^{\frac{1}{2}K_{aa}}, \qquad a = 1, \dots, K+M-1$$

Kulish, Sklianin'80-85

$$Q_{k,m}(u) = \widecheck{Q}_{k+m}(\widecheck{u}), \qquad \widecheck{u}_{i}^{(n)} = u_{i}^{(n)} - k + m$$

and Cartan matrix along the zigzag path

$$K_{ab} = (p_a + p_{a+1})\delta_{a,b} - p_{a+1}\delta_{a+1,b} - p_a\delta_{a,b+1}$$
 where 
$$p_n = \begin{cases} 1, & \text{if} \\ -1, & \text{if} \end{cases}$$

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# Conclusions and Prospects

- We proved Bazhanov-Reshetikhin formula for general fusion.
- We solved the associated Hirota discrete classical dynamics by generalized Baxter T-Q relations, found new Q-Q bilinear relations, reproduced nested TBA eqs. An alternative to the algebraic Bethe ansatz.
- Possible generalizations: noncompact irreps, mixed (covariant+contravariant) irreps, osp(n|2m) algebras.
   Trigonometric and elliptic(?) case.
- Non-standard R-matrices, like Hubbard or su(2|2) S-matrix in AdS/CFT, should be also described by Hirota eq. with different B.C.
- A potentially powerful tool for studying supersymmetric spin chains and 2d integrable field theories, including classical limits.

$$\frac{T(1,1,u+2)}{\phi(u+3)} = -\frac{Q_{1,0}(u+4)}{Q_{1,0}(u+2)}\frac{\phi(u+2)}{\phi(u+4)} + \frac{Q_{1,0}(u+4)}{Q_{1,0}(u+2)}\frac{Q_{1,1}(u)}{Q_{1,1}(u+2)}\frac{\phi(u+2)}{\phi(u+4)} - \frac{Q_{1,1}(u)}{Q_{1,1}(u+2)}$$