

Universality of Local Bulk Regime for Hermitian Matrix Models

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1 Introduction

1.1 Asymptotic "Philosophy" of RMT

Let M_n be a $n \times n$ S. O. matrix with i.i.d. entries $M_{ij} \sim \mathcal{N}(0, 1/n)$.

(ii) $\overline{N}_n ! N$ weakly as $n ! 1$;

(ii) $\overline{N}_n \ll N$ weakly as $n \rightarrow 1$;

(iii) $N(\cdot) = \int_{\mathbb{R}} (\cdot) d\mu$:

Property (ii) fixes the **global scale** of the spectral axis, yielding

$$\|f_l^{(n)}\|_2^2 : l = 1, \dots, n \approx nN(\cdot),$$

i.e., $(nN(\cdot))^{-1}$ is the typical eigenvalue spacing for large n in \mathbb{R} .

(ii) $\overline{N}_n \rightarrow N$ weakly as $n \rightarrow \infty$;

(iii) $N(\lambda) = \int_{\mathbb{R}} \chi_{[\lambda, \infty)}(x) d\mu(x)$;

Property (ii) fixes the **global scale** of the spectral axis, yielding

$$\int_{\mathbb{R}} f(x) d\mu(x) = \int_{\mathbb{R}} f(x) dN(x),$$

i.e., $(nN(\lambda))^{-1}$ is the typical eigenvalue spacing for large n in \mathbb{R} .

Assume $\text{supp } N$, then

λ does not depend on n (in the global scale), i.e., $nN(\lambda) \rightarrow n$: **global (macroscopic) regime**;

(ii) $\overline{N}_n \rightarrow N$ weakly as $n \rightarrow \infty$;

(iii) $N(\lambda) = \int_{\mathbb{R}} \rho(\lambda) d\lambda$;

Property (ii) fixes the **global scale** of the spectral axis, yielding

$$|f_l^{(n)}| \approx 2^{-l} : l = 1, \dots, n \approx nN(\lambda),$$

i.e., $(nN(\lambda))^{-1}$ is the typical eigenvalue spacing for large n in \mathbb{R} .

Assume $\text{supp } N$, then

$nN(\lambda) \gg 1$ does not depend on n (in the global scale), i.e., $nN(\lambda) \gg 1$: **global (macroscopic) regime**;

$nN(\lambda) \ll 1$: **local (microscopic) regime**;

(ii) $\overline{N}_n \rightarrow N$ weakly as $n \rightarrow \infty$;

(iii) $N(\lambda) = \int_{\mathbb{R}} \chi_{(\lambda, \infty)}(x) d\mu(x)$;

Property (ii) fixes the **global scale** of the spectral axis, yielding

$$\|f_l^{(n)}\|_2^2 \sim \int_{\mathbb{R}} \chi_{(\lambda, \infty)}(x) d\mu(x) \sim nN(\lambda),$$

i.e., $(nN(\lambda))^{-1}$ is the typical eigenvalue spacing for large n in \mathbb{R} .

Assume $\mu \ll \nu$ on $\text{supp } N$, then

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Probability theory analogs: LLN, CLT, "collective theorems", *Yu. Linnik*

"emerging universality" QFT

1.2 Linear Eigenvalue Statistics.

Take $\psi : \mathbb{R} \rightarrow \mathbb{R}$ and write the *linear eigenvalue statistic*

$$N_n[\psi] := \sum_{l=1}^n \psi \left(\lambda_l^{(n)} \right) = \text{Tr} \psi (M_n) = \int_{\mathbb{R}} \psi(x) N_n(dx):$$

ψ is known as the *test function*. In particular

$$\begin{aligned} N_n(\delta) &:= \sum_{l=1}^n \mathbb{1}_{\left[\lambda_l^{(n)} \in \left[\frac{1}{n}, \frac{2}{n} \right] \right]} ; l = 1, \dots, n \\ &= \sum_{l=1}^n \mathbb{1}_{\left[\lambda_l^{(n)} \in \left[\frac{1}{n}, \frac{2}{n} \right] \right]} = N_n[\delta] \end{aligned}$$

is the Eigenvalue Counting Measure of eigenvalues and $N_n = n^{-1} N_n$.

Define

$$\text{bulk } N = f \text{ } 2 \text{ supp } N : \vartheta > 0; \lim_{n \rightarrow \infty} \sup_j \inf_j j(\cdot) \text{ } n(\cdot) j = 0 g:$$

We have for $N_n[\cdot]$:

\cdot is n -independent: global regime;

$\cdot_n = \cdot((\cdot) L_n); L_n \rightarrow 1; n L_n \rightarrow 0$: intermediate bulk regime;

$\cdot_n = \cdot((\cdot) n \text{ } n(\cdot))$: local bulk regime

1.3 Typical Problems

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(ii) $\text{Var} f N_n g$, global and intermediate regime, fluctuations, "CLT"

1.3 Typical Problems

- (i) $\lim_{n \rightarrow \infty} \frac{1}{n} \sum_{i=1}^n N_n(\epsilon_i)$, global regime, selfaveraging, "LLN"
- (ii) $\text{Var} \sum_{i=1}^n N_n(\epsilon_i)$, global and intermediate regime, fluctuations, "CLT"
- (iii) $\mathbf{P} \sum_{i=1}^n N_n(\epsilon_i) = kg; k \in \mathbb{N}; E_n \sum_{i=1}^n N_n(\epsilon_i) = \mathbf{P} \sum_{i=1}^n N_n(\epsilon_i) = 0g$,
 gap probability, local regime, spacings, universality
 in particular $\sum_{i=1}^n N_n(\epsilon_i) = \sum_{i=1}^n N_n(\epsilon_i + s)$; $s \in \mathbb{N}$ bulk N for the
 local bulk regime

1.4 Hermitian Matrix Models

$n \times n$ hermitian random matrices with the law

$$P_n(dM) = Z_n^{-1} \exp \left(-n \operatorname{Tr} V(M) \right) dM;$$

$$dM = \prod_{j=1}^n dM_{jj} \prod_{1 \leq j < k \leq n} dM_{jk} dM_{kj};$$

$V : \mathbb{R} \rightarrow \mathbb{R}_+$ is a continuous function (potential), and

$$V''(x) > 0; \quad L < 1 \quad V(x) \sim (2 + \epsilon) \log(1 + |x|) > 0; \quad |x| \leq L$$

$V = x^2/2$ corresponds to the Gaussian Unitary Ensemble (GUE).

1.5 Results (a collection)

(i) For any probability measure m on \mathbb{R} ; $m(\mathbb{R}) = 1$ define (Gauss)

$$E[m] = \int V(\cdot) m(d) \quad \int \int \log j \quad j m(d) m(d);$$

and let N

In particular

$$\text{v.p.} \int_{\text{supp } N} \frac{(\cdot) d}{Z} = V^0(\cdot) = 2; \quad 2 \text{ supp } N:$$

i.e., an analog of the LLN : *Wigner, 52; Brezin et al, 79; A. Boutet de Monvel, P., Shcherbina, 95; Deift et al 98; Johansson, 98; P., Shcherbina, 07.*

(ii) $\text{Var } fN_n[\cdot]g$ does not grow with n if $\cdot = 2$

In particular

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V^0 is Lip 1;

there exists a closed interval $[a; b] = \text{supp } N$ such that

$$\sup_{2[a;b]} |jV^{000}(\cdot)| \leq C_1 < 1; \quad 0 < \inf_{2[a;b]} (\cdot):$$

Then we have for any $d > 0$:

(i)

$$\sup_{2[a+d;b-d]} |j_n(\cdot)| \leq Cn^{-2=9};$$

i.e., $[a + d; b - d]$ belongs to bulk N ;

(ii) if $p_l^{(n)}$; $l = 1; 2; \dots$ are the marginals of the joint probability density of eigenvalues, then for any $\alpha \in [a + d; b - d]$

$$\lim_{n \rightarrow \infty} [p_l^{(n)}] = p_l^{(\infty)} = \alpha + \frac{x_l}{n - \alpha}; \dots; \alpha + \frac{x_l}{n - \alpha}$$

(iii) if $E_n(\lambda) = \mathbf{P} f_{\lambda}^{(n)} \geq \lambda; l = 1; \dots; ng$ is the gap probability of the ensemble and $\lambda = (0; 0 + s = n(0)) \in [a + d; b - d]$, then

$$\lim_{n \rightarrow \infty} \mathbf{P}(\lambda) = \det(1 - S(s));$$

where

$$(S(s)f)(x) = \int_0^s \frac{\sin(x-y)}{(x-y)} f(y) dy; x \in [0; s];$$

Dyson, 61, 73; P. Shcherbina, 97, 07; Deift et al 99

More if $\lambda = 0; 1$ (singular points (e.g. edge) universality).

2 Proof (outline)

2.1 Orthogonal Polynomials Techniques

Weyl integration formula for the joint eigenvalue density

$$\rho_n(\lambda_1, \dots, \lambda_n) = Q_n^{-1} e^{-\sum_{k=1}^n V(\lambda_k)}$$

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Weyl integration formula for the joint eigenvalue density

$$\rho_n$$

$$P_l^{(n)} = e^{-nV/2} P_l^{(n)} \text{ and } K_n(\lambda; \mu) = \prod_{l=0}^{n-1} P_l^{(n)}(\lambda) P_l^{(n)}(\mu),$$

$$K_n(\lambda; \mu) K_n(\mu; \nu) d\mu = K_n(\lambda; \nu):$$

$$I^{(n)} = e^{-nV/2} P_I^{(n)} \text{ and } K_n(\cdot; \cdot) = \prod_{l=0}^{n-1} I^{(n)}(\cdot) I^{(n)}(\cdot),$$

$$K_n(\cdot; \cdot) K_n(\cdot; \cdot) d = K_n(\cdot; \cdot):$$

$$p_i^{(n)} = e^{-nV} \int \prod_{l=0}^{n-1} p_l^{(n)} \text{ and } K_n(\mathbf{d}) = \int \prod_{l=0}^{n-1} p_l^{(n)}(\mathbf{d}_l) p_l^{(n)}(\mathbf{d}_l),$$

$$K_n(\mathbf{d}) K_n(\mathbf{d}') d = K_n(\mathbf{d}; \mathbf{d}')$$

Then the marginals $p_i^{(n)}$ of p_n are given by the **determinant formulas**

$$\begin{aligned} p_i^{(n)}(\mathbf{d}_1, \dots, \mathbf{d}_i) &= \int p_n(\mathbf{d}_1, \dots, \mathbf{d}_n) d_{i+1} \dots d_n \\ &= (n \dots (n-i+1))^{-1} \det f K_n(\mathbf{d}_j; \mathbf{d}_k) g_{j,k=1}^i \end{aligned}$$

Determinant formulas imply:

$$(a) \int f n^{-1} N_n[\mathbf{g}] g = \int \prod_{l=0}^{n-1} p_l^{(n)}(\mathbf{d}_l) d_l; \quad p_n(\mathbf{d}) = K_n(\mathbf{d}; \mathbf{d}')$$

$$(b) \text{Var} fN_n[\cdot]g = \frac{1}{2} \int \int_{\mathbb{R} \times \mathbb{R}} (f(x_1) - f(x_2))^2 K_n^2(x_1; x_2) d x_1 d x_2;$$

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(c) $\mathbf{P}fN_n(\cdot) = 0g = \det(1 - K_n(\cdot)),$ where

$$(K_n(\cdot)f)(x) = \int_{\mathbb{Z}} K_n(x; y)f(y)$$

These asymptotics were found by *Deift et al, 99* and lead to the solution of (iii), i.e., the universality of the local bulk regime of the hermitian matrix models, and (ii), i.e., the limiting laws of fluctuations of linear eigenvalue statistics, the CLT in particular, *P, 06*, for real analytic potentials.

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On the other hand, in *P., Shcherbina 97, 07* the universality of the local bulk regime of hermitian matrix models is proved for globally C^2 and locally C^3 potentials (see above theorem), basing on the orthogonal polynomial techniques, in particular on the above integral representation for K_n , but **NOT** using asymptotics of corresponding orthogonal polynomials.

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In *P., Shcherbina, 97* $\sin(x) = (x)$ is obtained via its Taylor expansion.

In *P., Shcherbina, 07* $\sin(x) = (x)$ is obtained as solution of a non-linear integro-differential equation.

2.2 Integro-differential Equation for Rescaled Reproducing Kernel

We start from the integral representation à la determinant formulas

$$n^{-1} K_n(x; y) = Q_{n;2}^{-1} e^{-n(V(x)+V(y))} \int_{\mathcal{Z}} \prod_{j=2}^n d_j e^{nV(x_j)} \prod_{j=2}^n \prod_{2 \leq j < k \leq n} (x_j - x_k)^2$$

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$$\frac{\partial}{\partial x} K_n(x; y) = \frac{1}{2} V^\ell(0 + x=n) K_n(x; y) + \int \frac{K_n(x^\ell; x^\ell) K_n(x; y) - K_n(x; x^\ell) K_n(x^\ell; y)}{x - x^\ell} dx^\ell.$$

Differentiate the representation with respect to x to obtain the identity

$$\frac{\partial}{\partial x} K_n(x; y) = \frac{1}{2} V^{\ell}(0 + x=n) K_n(x; y) + \sum K_n(x^{\ell}; x^{\ell}) K_n(x; y) - K_n(x; x^{\ell})$$

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and write

$$\frac{\partial}{\partial X} K_n(x; y) = \int \frac{K_n(x; x^\theta) K_n(x^\theta; y)}{Z}$$

and write

$$\frac{\partial}{\partial X} K_n(x; y) = \int \frac{K_n(x; x') K_n(x'; y)}{Z}$$

We prove next that under the conditions of theorem we have uniformly in

$|x|, |y| < L; \quad 0 \leq [a + d; b - d]$:

$$\frac{\partial}{\partial x} K_n(x; y) + \frac{\partial}{\partial y} K_n(x; y) = C n^{-1/8} + |x - y| n^{-2} ;$$

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$$|K_n(x; y) - K_n(0; y - x)| \leq C|x| n^{-1/8} + |x - y| n^{-2};$$

$$\frac{\partial}{\partial x} K_n(x; y) \leq C \int_{|x|}^L dx \left(\frac{\partial}{\partial x} K_n(x; y) \right)^2 \leq C;$$

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$$|K_n(x; y) - K_n(0; y - x)| \leq C|x| n^{-1/8} + |x - y| n^{-2};$$

$$\frac{\partial}{\partial x} K_n(x; y) \leq C \int_{|x|}^L dx \left(\frac{\partial}{\partial x} K_n(x; y) \right)^2 \leq C;$$

Now, if

$$K_n(x) = K_n(x;0) \mathbf{1}_{|x| \leq L} + K_n(L;0)(1 + L - x) \mathbf{1}_{L < x \leq L+1} \\ + K_n(-L;0)(1 + L + x) \mathbf{1}_{-L-1 \leq x < -L};$$

then for $|y| \leq L$

$$\frac{\partial}{\partial y} K_n(y) = \int_{|x'| \leq 2L-3} \frac{K_n(x') K_n(y - x')}{x'} dx' + O(L^{-1});$$

and $\int_{|x| \leq L} K_n(x)^2 dx \leq 1$; $\int_{|x| \leq L} \frac{d}{dx} K_n(x)^2 dx \leq 1$;

2.3 Asymptotic Solution of Equation

Consider the Fourier transform

$$K_n(p) = \int K_n(x) e^{ipx} dx; \quad K_n(x) = (2\pi)^{-1} \int K_n(p) e^{-ipy} dp:$$

Then we have from $n^{-1} K_n(\cdot; \cdot) =$

Since K_n is "asymptotically even"

$$\int_{-\infty}^{\infty} jK_n(p) \int_{-\infty}^{\infty} K_n(p) j^2 dp$$

$$= 2 \int_{-\infty}^{\infty} jK_n(x) \int_{-\infty}^{\infty} K_n(x) j^2 dx \quad Cn^{-1/8} \log^3 n:$$

we obtain the Fourier form of the above integro-differential equation:

$$\int_{-\infty}^{\infty} K_n(p) \int_0^{\infty} K_n(p^l) dp^l \int_{-\infty}^{\infty} p e^{ipy} dp = O(L^{-1}); \quad |y| \leq L=3:$$

Since K_n is "asymptotically even"

$$\int_{\mathbb{Z}} j\hat{K}_n(p) \hat{K}_n(p) j^2 dp = 2 \int_{\mathbb{Z}} jK_n(x) K_n(x) j^2 dx = C n^{-1/8} \log^3 n;$$

we obtain the Fourier form of the above integro-differential equation:

$$\int_{\mathbb{Z}} \hat{K}_n(p) \int_0^{\mathbb{Z}} \hat{K}_n(p^l) dp^l \int_{\mathbb{Z}} p e^{ipy} dp = O(L^{-1}); \quad |jy| \leq L=3;$$

Besides, since K_n is positive definite \hat{K}_n is "asymptotically non-negative":

$$\int_{\mathbb{Z}} \hat{K}_n(p) j\hat{f}(p) j^2 dp = C \|j\hat{f}\|_{L^2(\mathbb{R})}^2 (n^{-1/8} \log^4 n + O(L^{-1}));$$

Now consider the functions

$$F_n(p) = \int_0^p k_n(p^\ell) dp^\ell:$$

Since $p k_n \in L^2(\mathbb{R})$, the sequence $\{F_n\}$ consists of functions that are of uniformly bounded variation, uniformly bounded and equicontinuous on \mathbb{R} .

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Thus $\{F_n\}$ is a compact family with respect to the uniform convergence.

Hence, the limit F of any subsequence $\{F_{n_k}\}$ possesses the properties:

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(b) $F(p) = F(-p)$;

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Hence, the limit F of any subsequence $\{F_{n_k}\}$ possesses the properties:

(a) F is bounded and continuous;

(b) $F(p) = F(p^-)$;

(c) $F(p) = F(p^\ell)$, if $p = p^\ell$;

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(e) F satisfies the following equation, valid for any smooth function g of compact support:

$$\int (F(p) - p)g(p) dF(p) = 0:$$

The last property implies that $F(p) = p$ or $F(p) = \text{const}$, hence it follows from (a) – (c) that

$$F(p) = p \mathbf{1}_{|p| \leq p_0} + \text{sign}(p) \mathbf{1}_{|p| > p_0};$$

where $p_0 = (0)$ by (d).

We conclude that the equation is uniquely soluble, thus the sequence $fF_n g$ converges uniformly on any compact to the above F . This implies the weak

convergence of the sequence $fK_n g$ to the function

$$K(x) = \frac{\sin\left(\frac{1}{n}x\right)}{\frac{1}{n}x};$$

convergence of the sequence $fK_n g$ to the function

$$K(x) = \frac{\sin(x)}{x};$$

But weak convergence implies

$$\lim_{n \rightarrow \infty} K_n(x; y) = K(x; y):$$

uniformly in $(x; y)$, varying on a compact set of \mathbb{R}^2 , because

$$\frac{d}{dx} K_n \in L^2(\mathbb{R}):$$